

Ultra-sensitive atomic spin measurements with a nonlinear interferometer

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NG roup
itchell

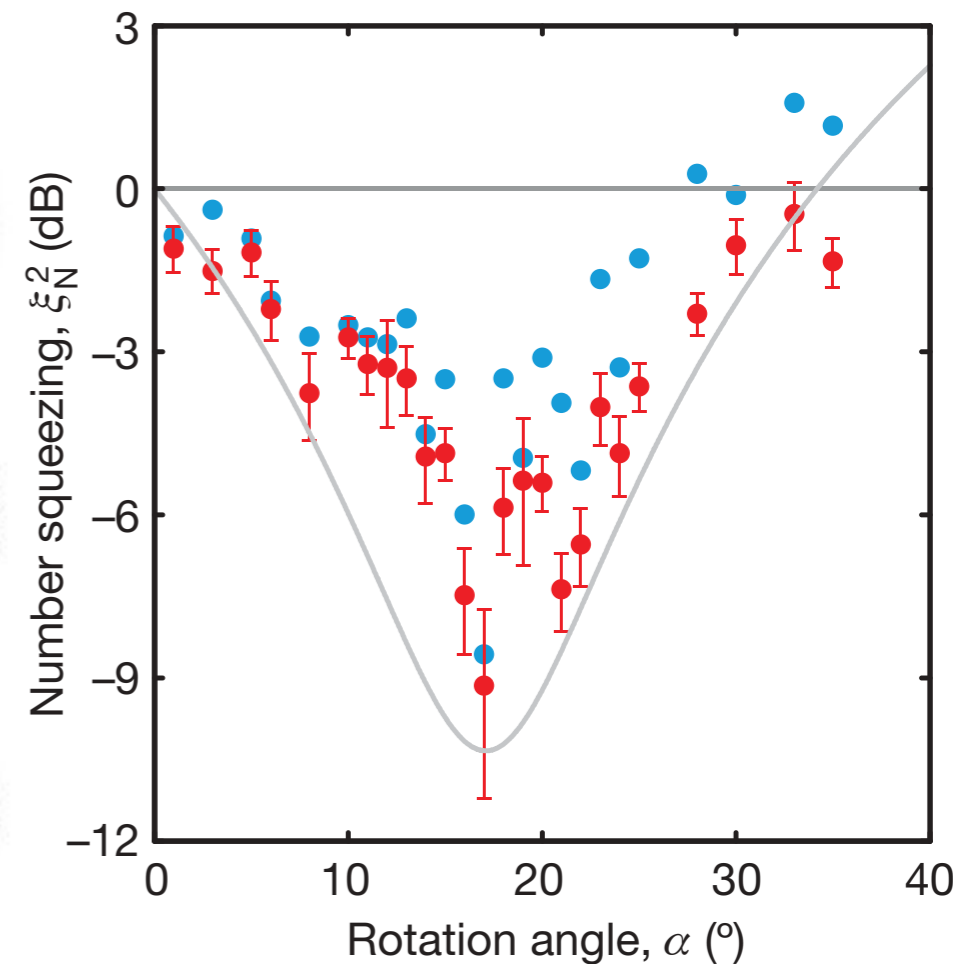
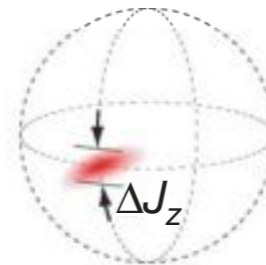
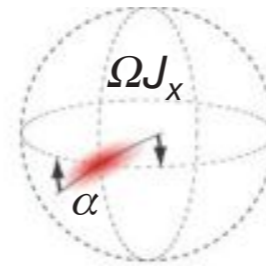
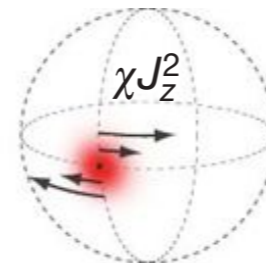
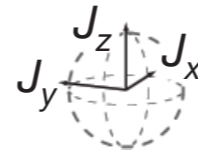
Nonlinear interferometers

Matter-wave interferometer

atom-atom interactions
lead to nonlinear rotations
& spin squeezing

Gross, Nature 464, 1165 (2010)

a



Nonlinear interferometers

Nonlinear faraday rotation
& optical magnetometry

nonlinear medium (atomic
gas) gives rise to an
enhanced rotation signal

Wojciechowski, PRA 81,
053420 (2010)

c.f. Budker group, Berkeley

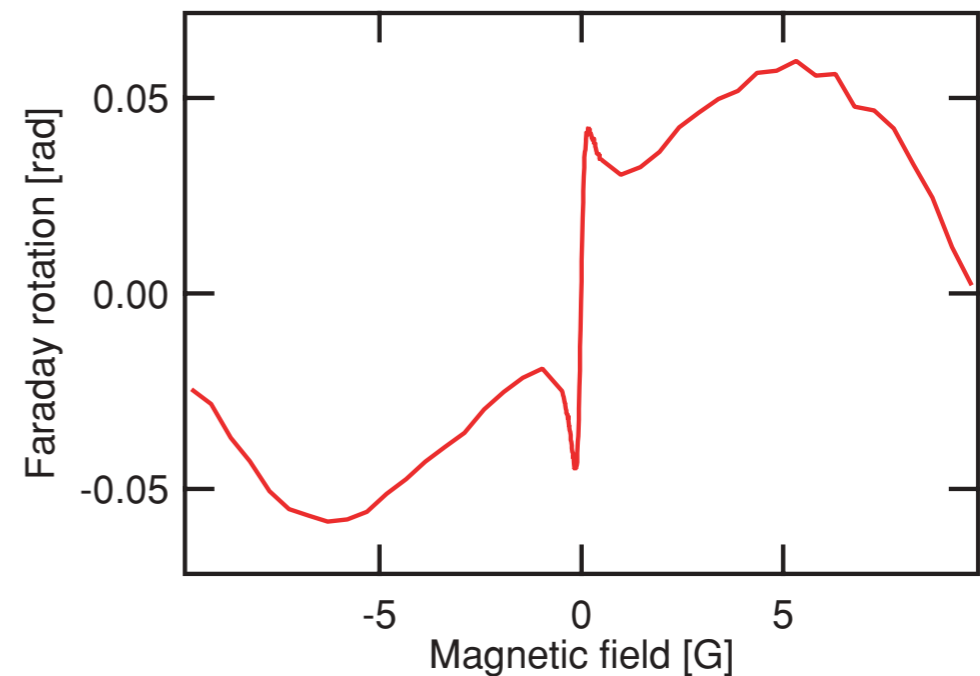
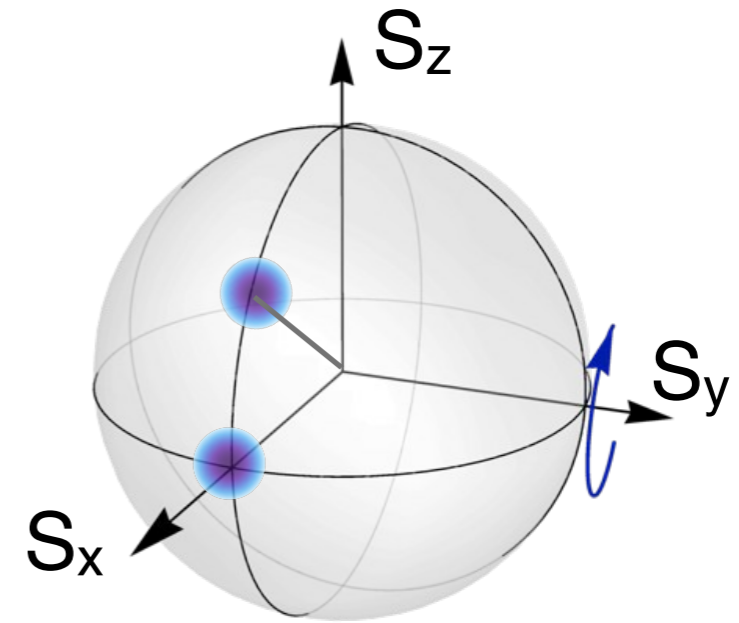
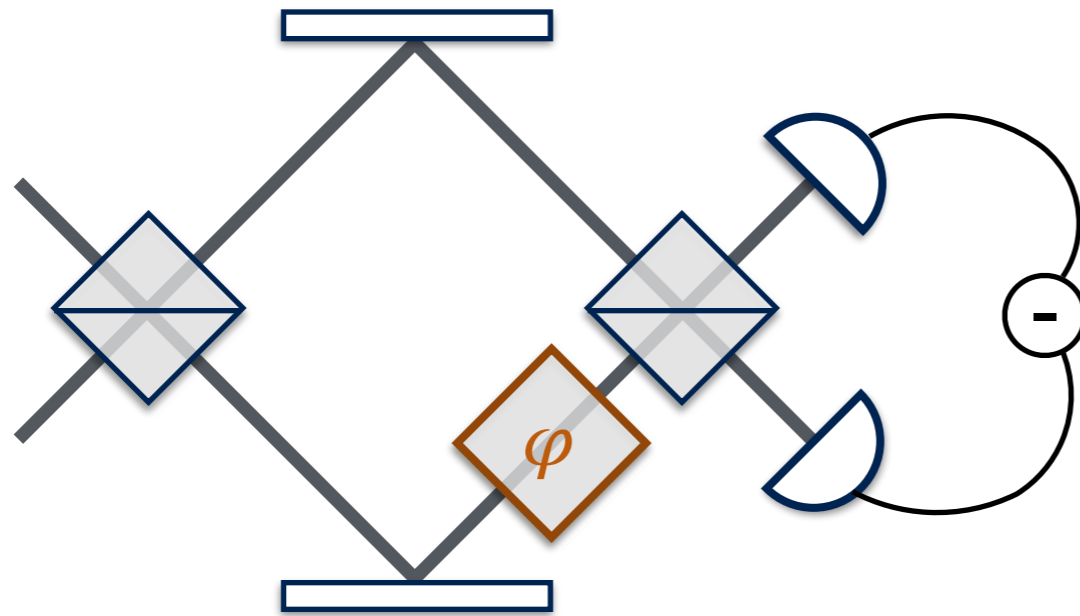


FIG. 2. (Color online) Linear (wide) and nonlinear (narrow) Faraday rotation resonances centered at $B = 0$. Signals were recorded at the time $\tau = 2$ ms after switching on the probing beam. The probe power is $64 \mu\text{W}$. At that power the NFR resonance is substantially power broadened but is well visible in comparison with the LFR. The slope of the central part is ≈ 0.6 rad/G.

Linear interferometer



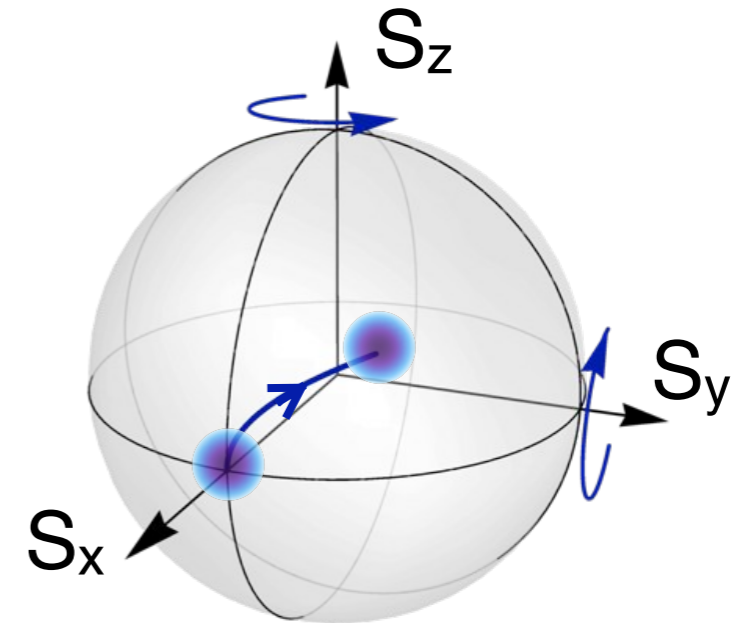
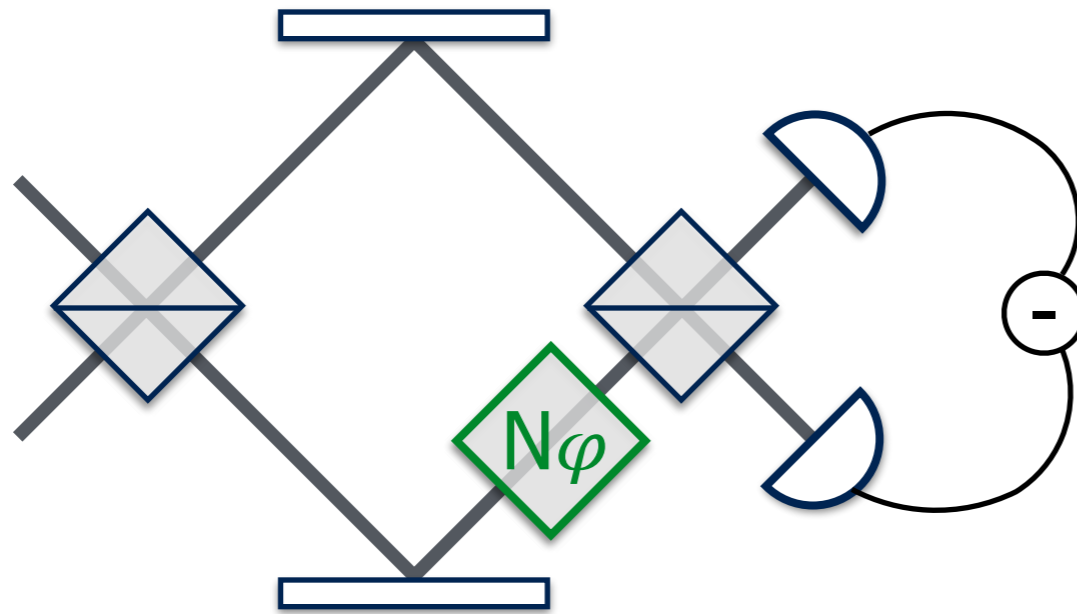
phase shift $\phi = \alpha J_y$

shot-noise $\Delta\phi = 1/\sqrt{N}$

sensitivity

$$\Delta J_y = \frac{\Delta\phi}{|\partial\phi/\partial J_y|} = \frac{1}{\alpha\sqrt{N}}$$

Nonlinear interferometer



phase shift $\phi = \beta N J_y$

shot-noise $\Delta\phi = 1/\sqrt{N}$

sensitivity

$$\Delta J_y = \frac{\Delta\phi}{|\partial\phi/\partial J_y|} = \frac{1}{\beta N^{3/2}}$$

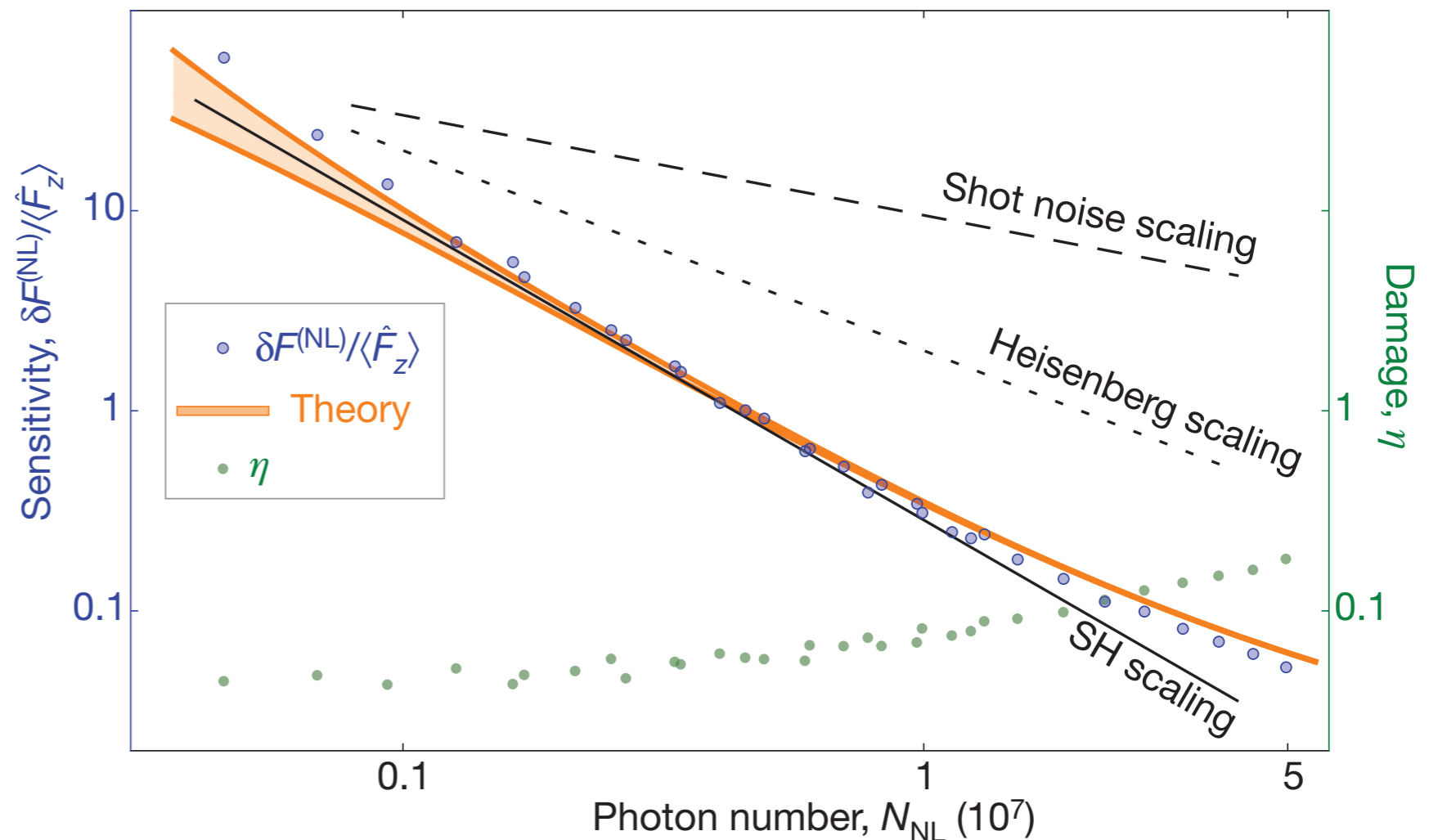
Can improved scaling give better sensitivity?

observed
improved scaling

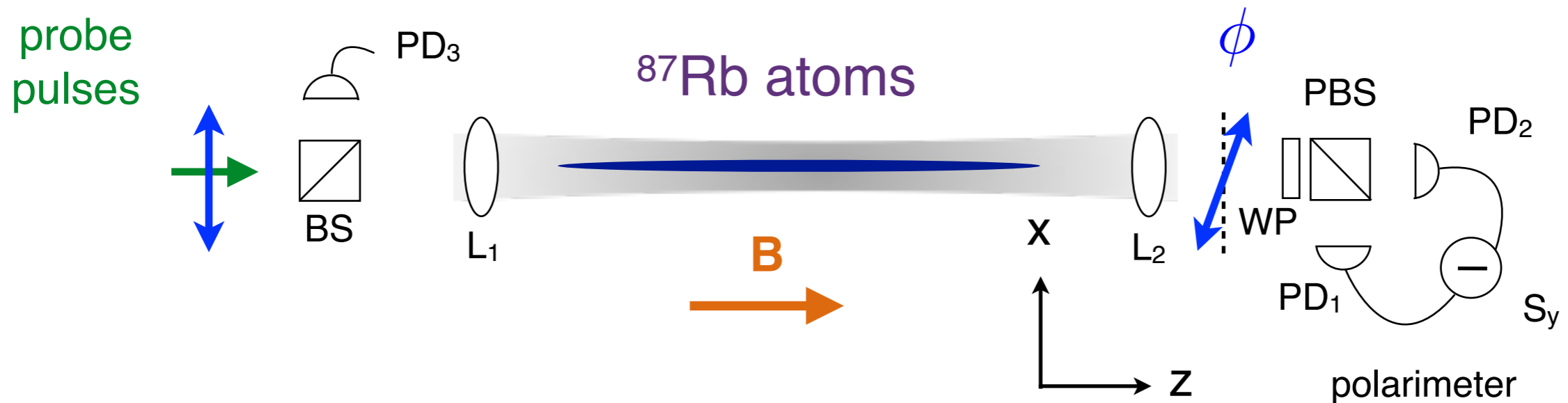
$$\Delta J \propto N^{-3/2}$$

did not continue to
large N, nor give
improved absolute
sensitivity

Napolitano, Nature
486, 471 (2011)



Quantum atom-light interface



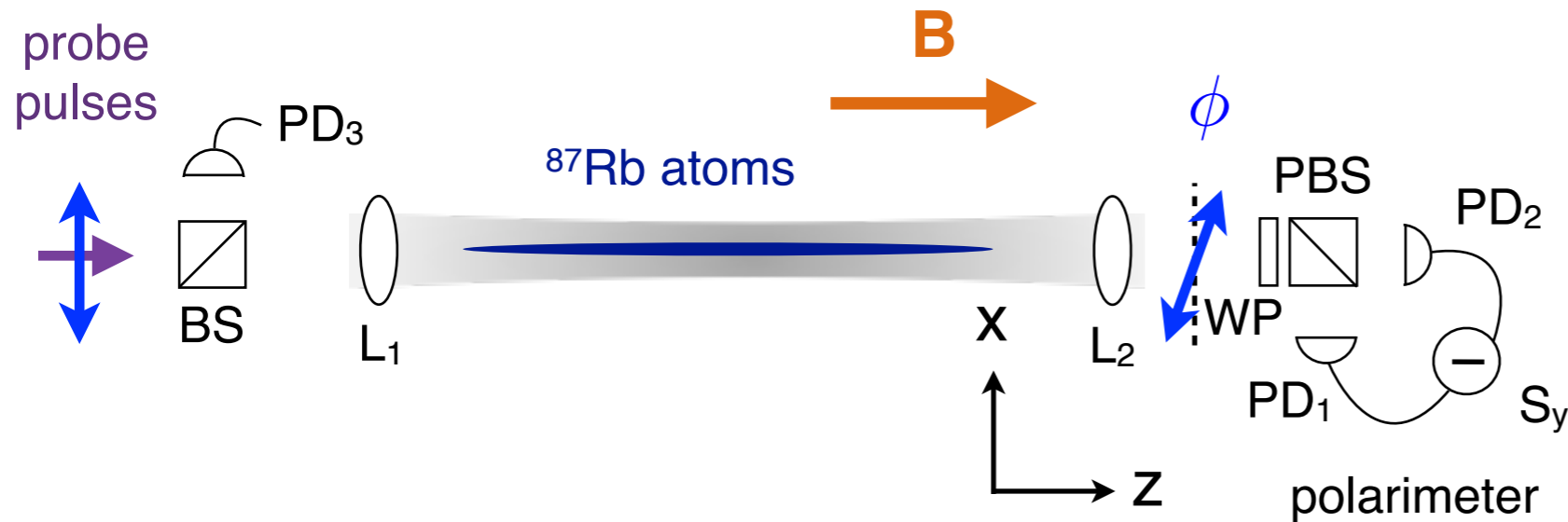
$\sim 10^6$ ^{87}Rb atoms at $25\mu\text{K}$
 $f=1$ ground-state

$1\mu\text{s}$ long pulses
linearly polarized
“mode matched” to atoms
 0.7 GHz from D_2 line

- ¹ effective OD > 50
- ² Sensitivity 512 spins, $< \text{SQL}$
- ³ QND measurement
- ⁴ spin squeezing

- 1 Kubasik, et al. PRA 79, 043815 (2009)
- 2 Koschorreck, et al. PRL 104, 093602 (2010)
- 3 Koschorreck, et al. PRL 105, 093602(2010)
- 4 Sewell, et al. PRL 109, 253605 (2012)

Atom-light interaction



$$\tau H = G_1 S_z J_z + G_2 S_x J_x + \gamma B_z J_z$$

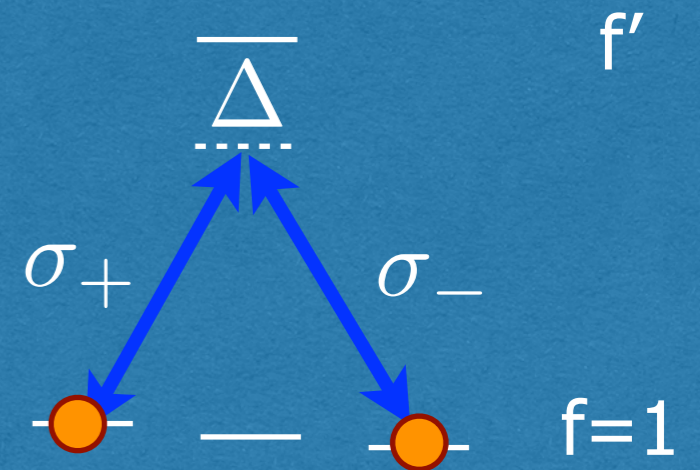
collective spin

$$\mathbf{J} = \sum_{i=1}^{N_A} \mathbf{j}^{(i)}$$

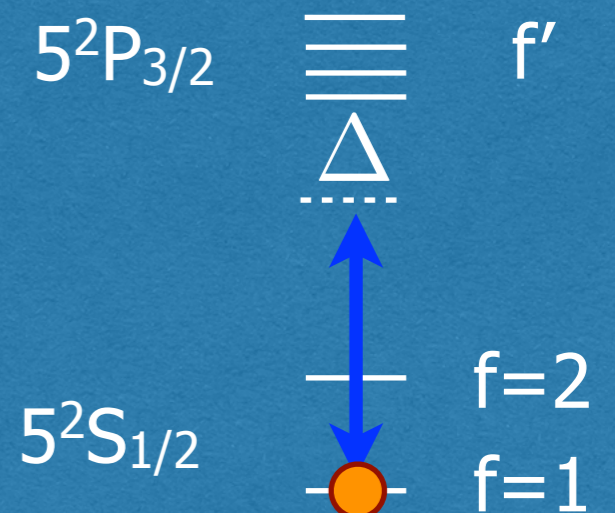
Stokes operators

$$\mathbf{S} \equiv \sum_{i=1}^{N_L} \mathbf{s}^{(i)}$$

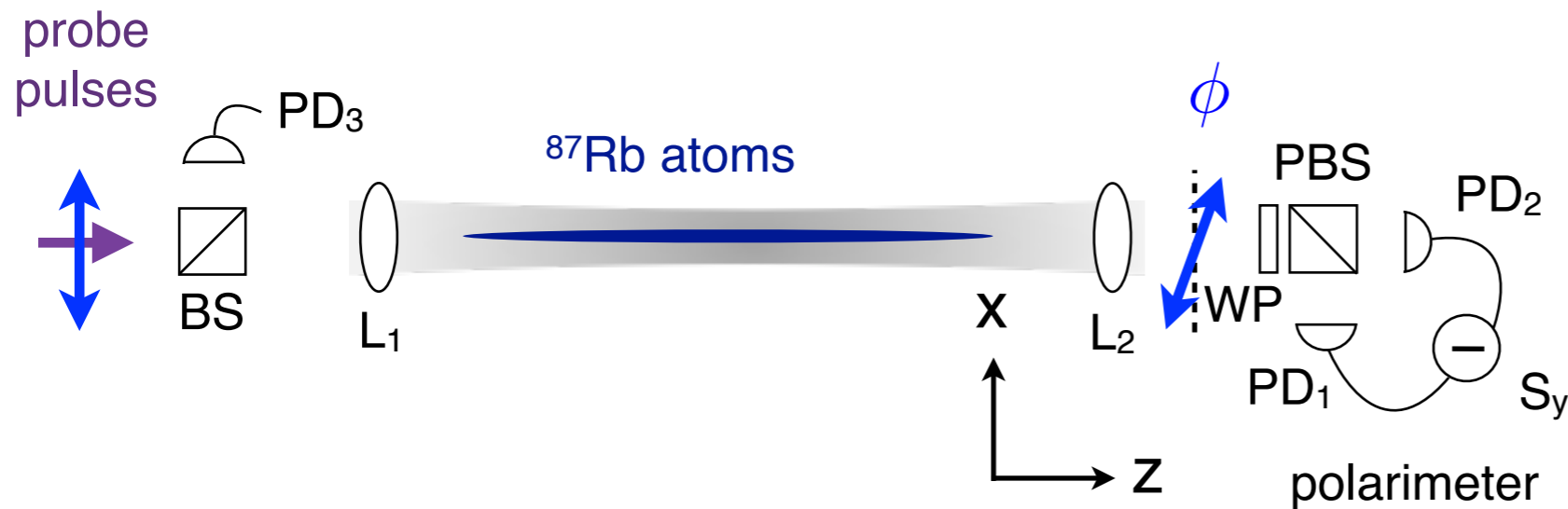
pseudo-spin



^{87}Rb D₂ line



Atom-light interaction



$$\tau H = G_1 S_z J_z + G_2 S_x J_x + \gamma B_z J_z$$

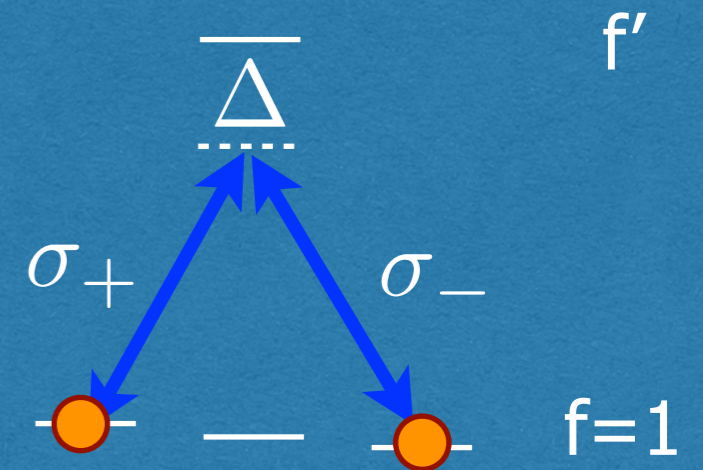
Faraday rotation

$$\phi = G_1 S_x J_z$$

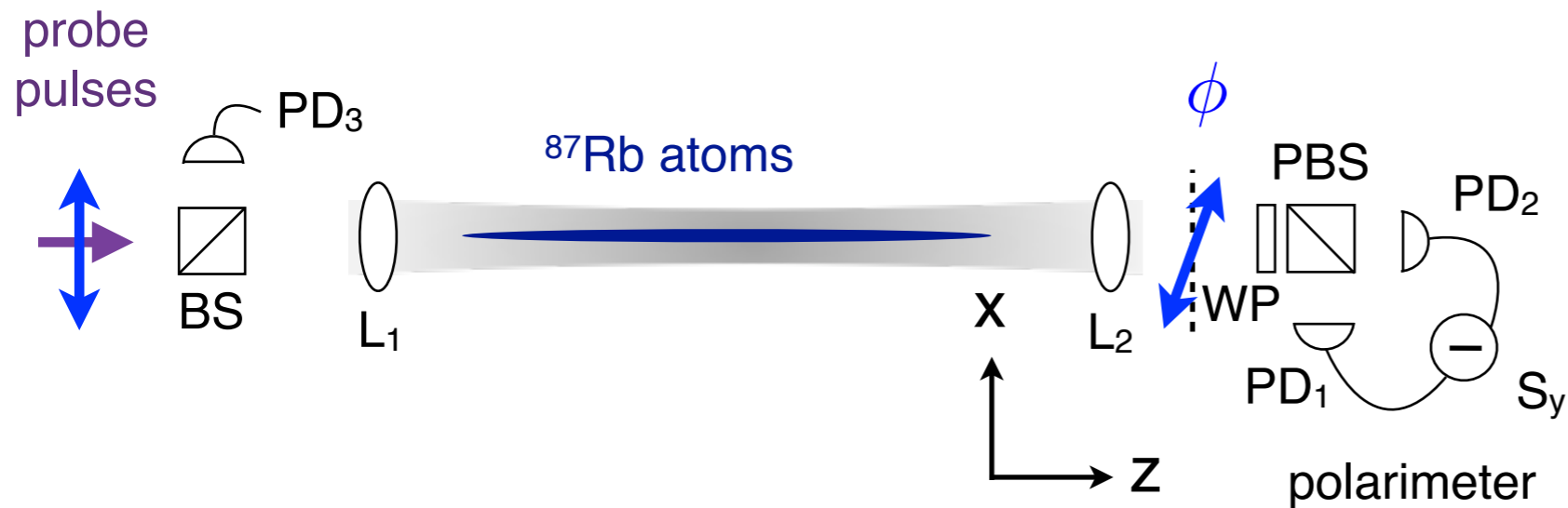
Larmor precession

$$J_y = \gamma B_z J_x$$

pseudo-spin



Atom-light interaction

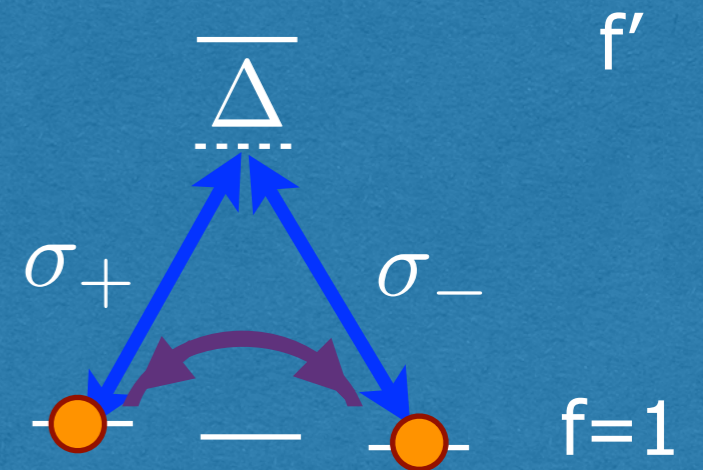


$$\tau H = G_1 S_z J_z + G_2 S_x J_x + \gamma B_z J_z$$

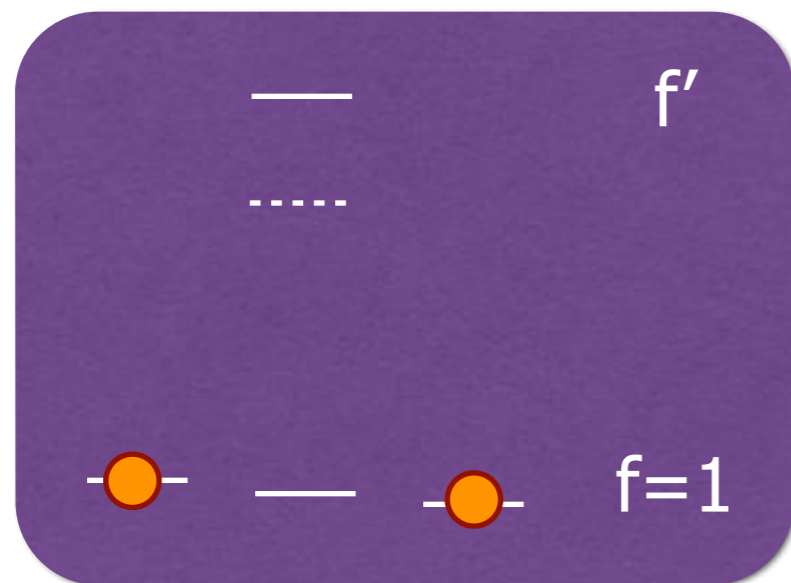
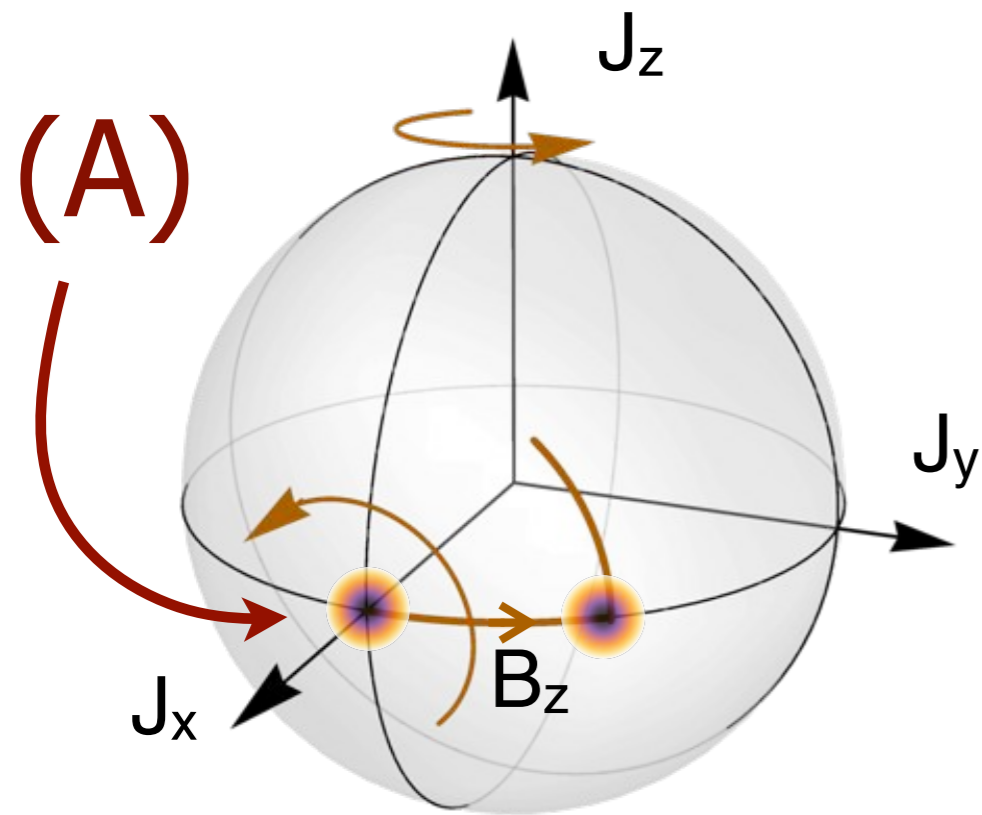
atomic birefringence

Raman transitions

pseudo-spin



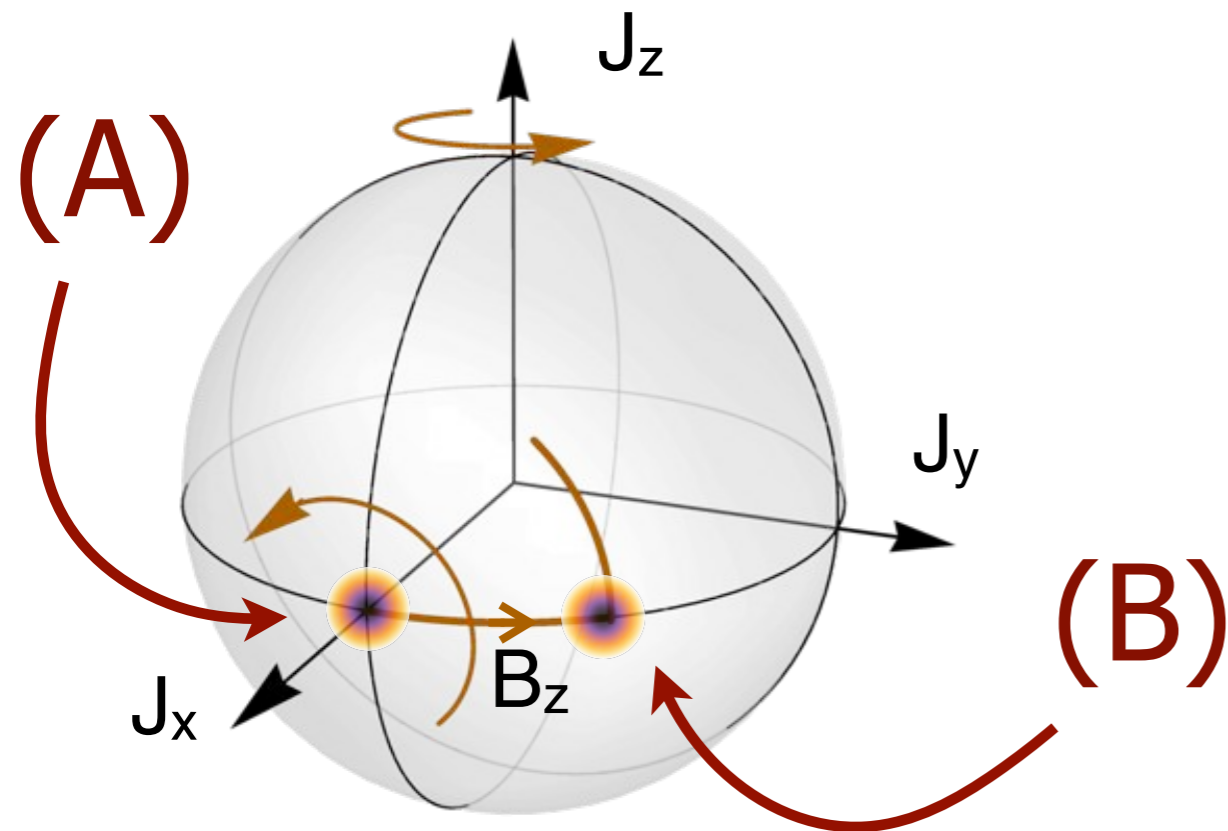
Atomic state



(A)

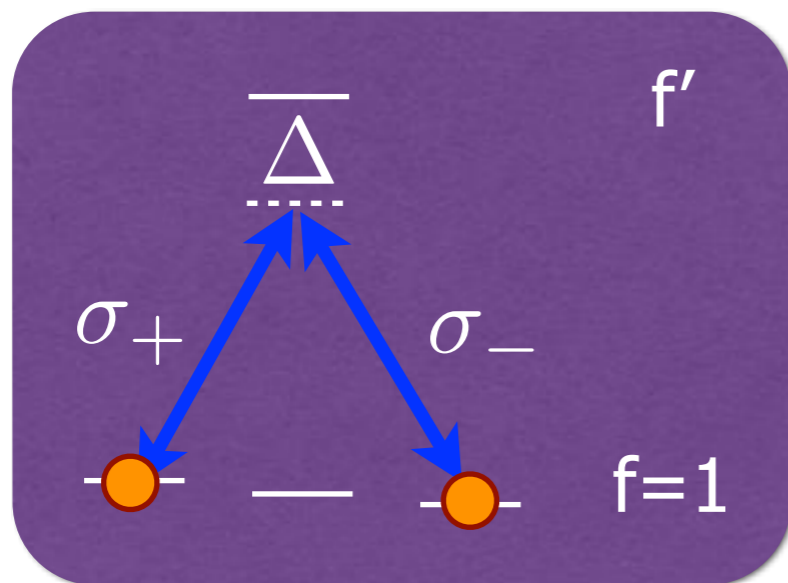
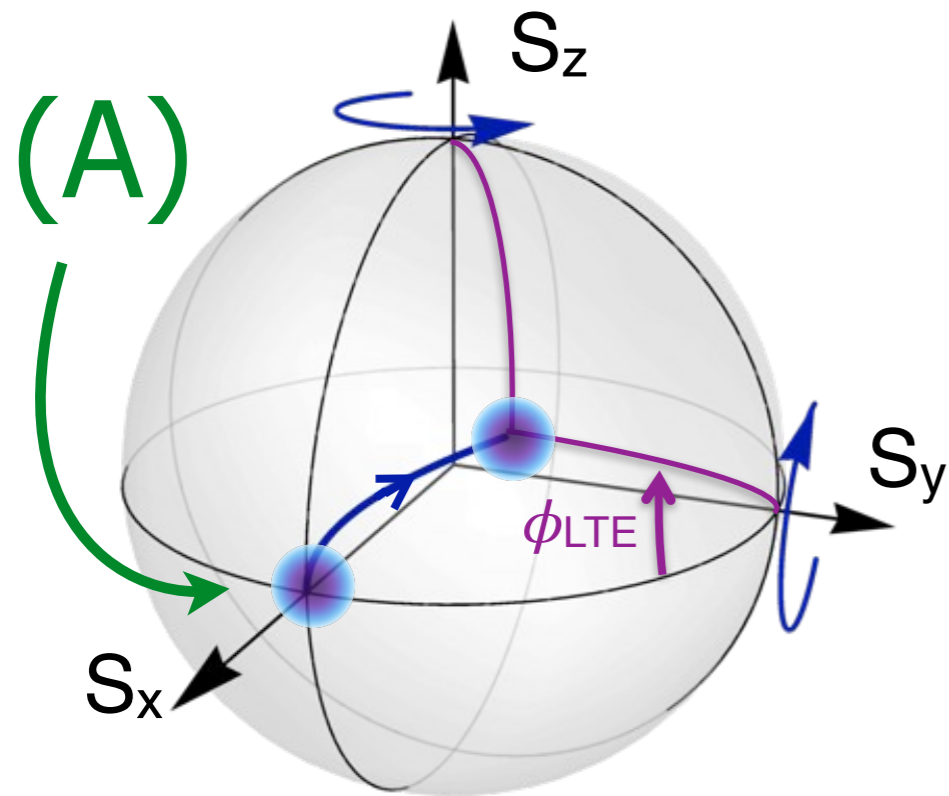
input $J_x = N_A/2$
aligned state
prepared via optical
pumping

Atomic state rotates due to B-field



- (A) input $J_x = N_A/2$
aligned state
prepared via optical
pumping
- (B) B_z rotates collective
atomic spin
 $J_x \Rightarrow J_y \propto B_z$

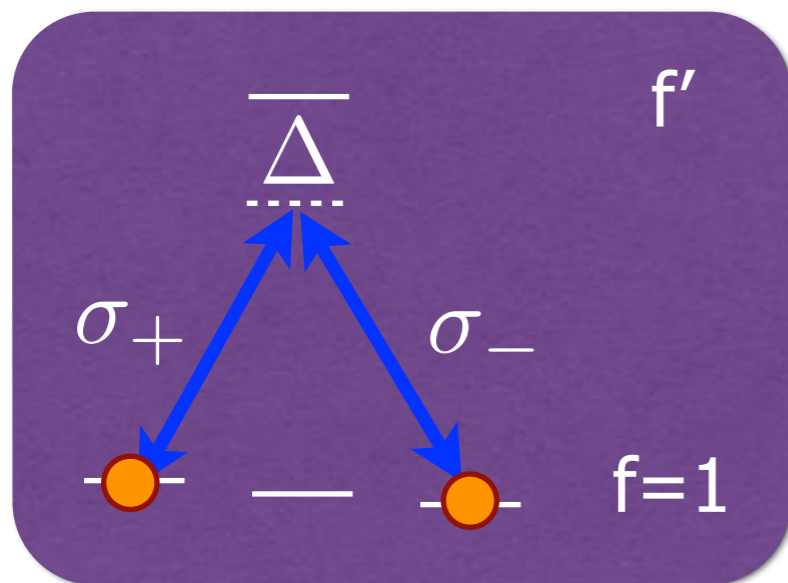
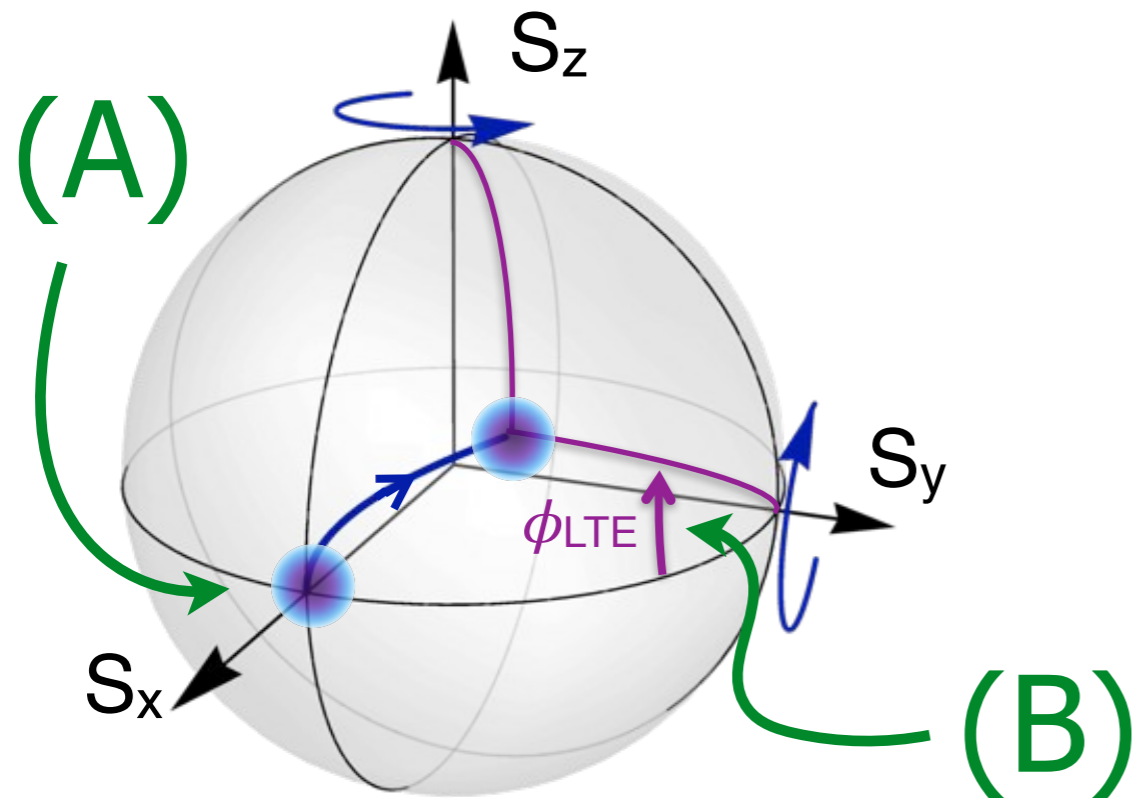
Read-out via optical birefringence



(A)

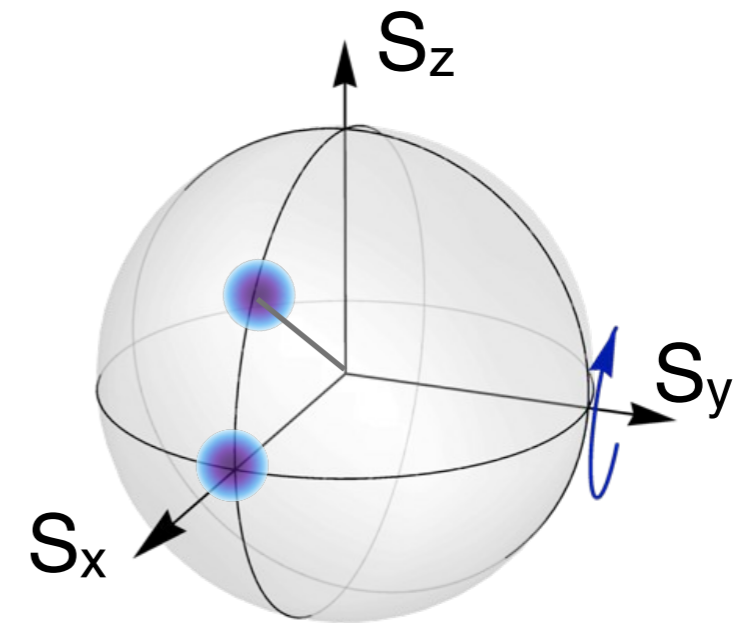
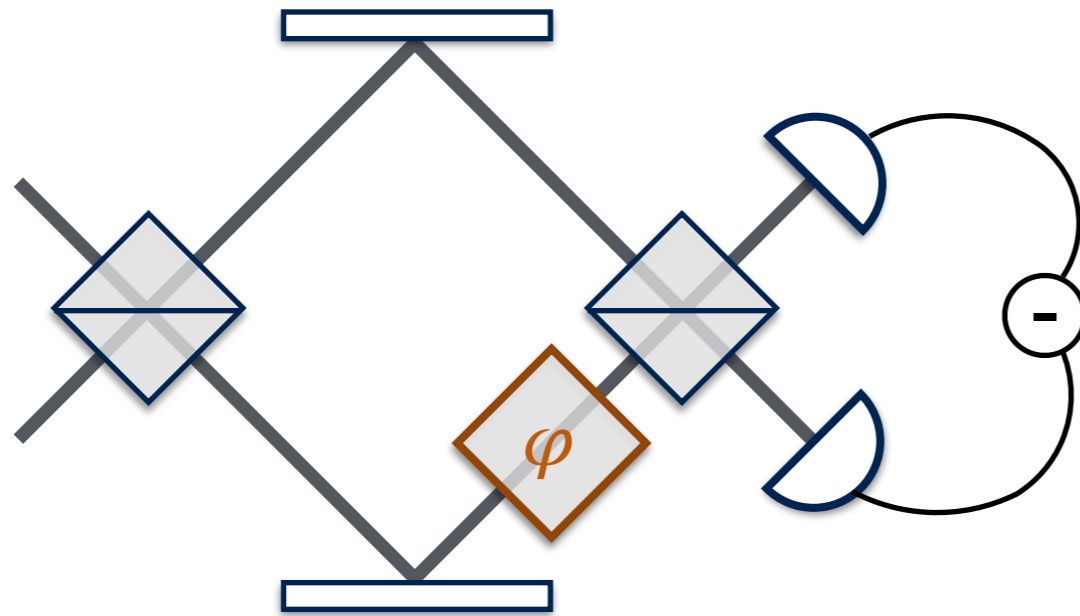
input $S_x = N_L/2$
polarised optical
pulse

Read-out via optical birefringence



- (A) input $S_x = N_L/2$
polarised optical pulse
- (B) rotates $S_x \Rightarrow S_z \propto J_y$
(atomic birefringence)

Linear interferometer



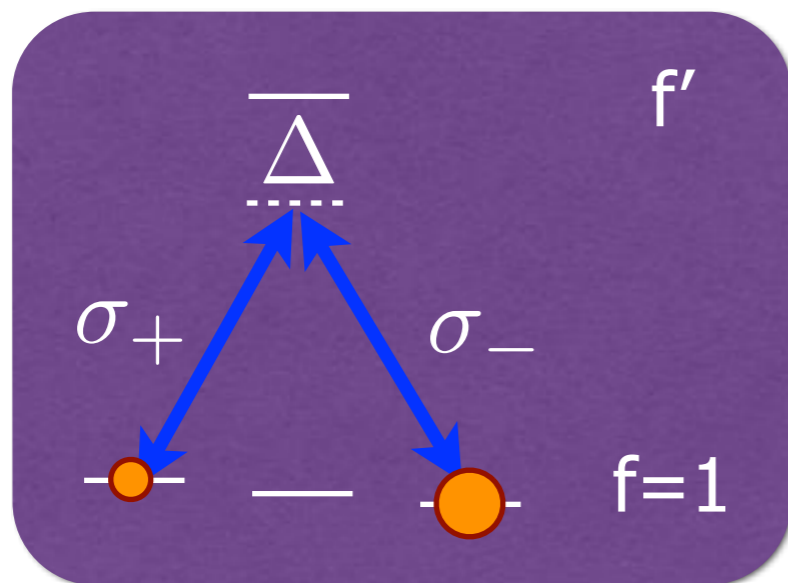
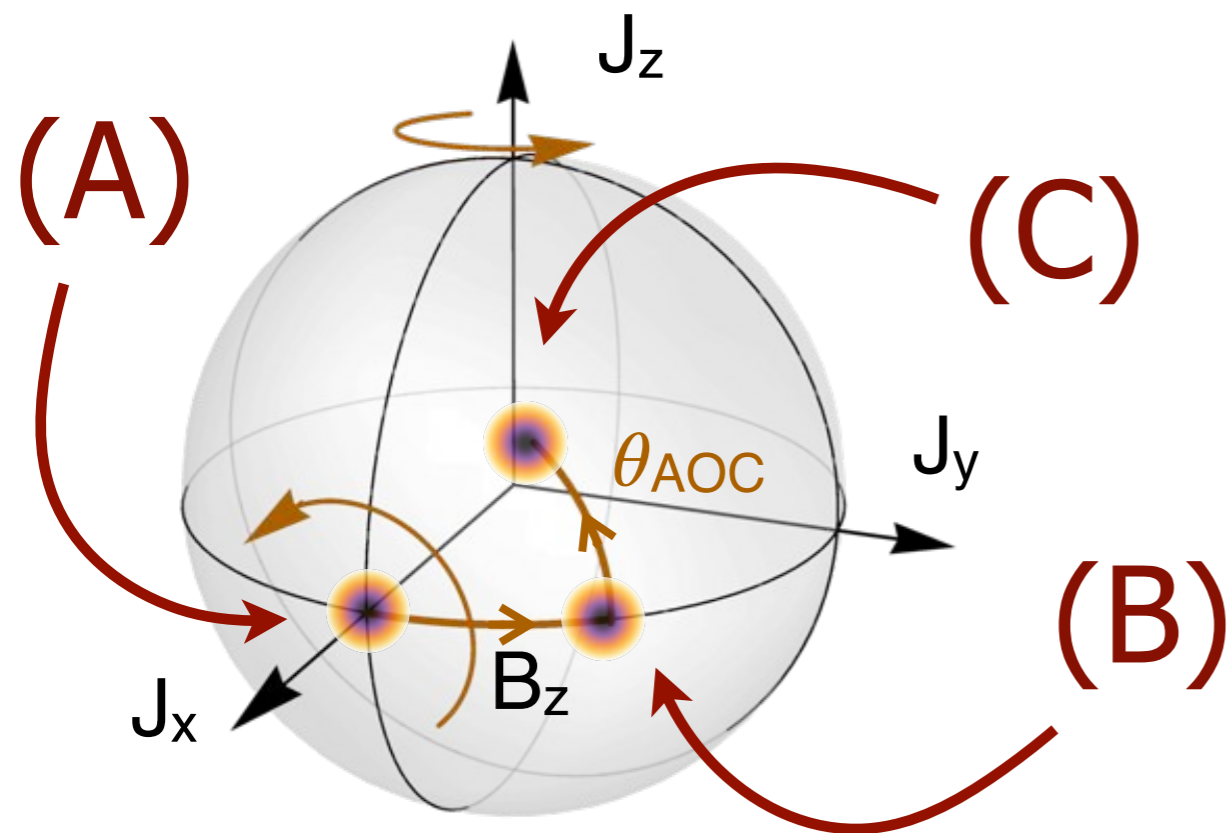
phase shift $\phi = G_2 J_y$

shot-noise $\Delta\phi = 4/\sqrt{N}$

sensitivity

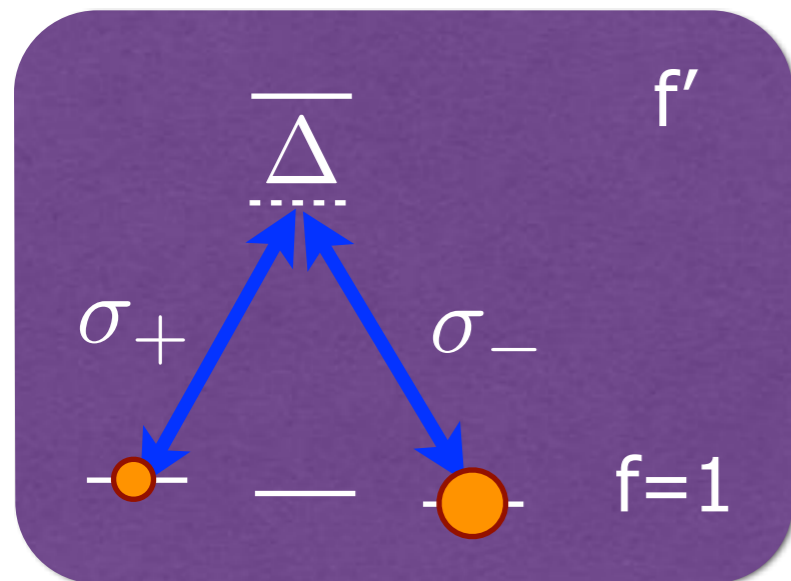
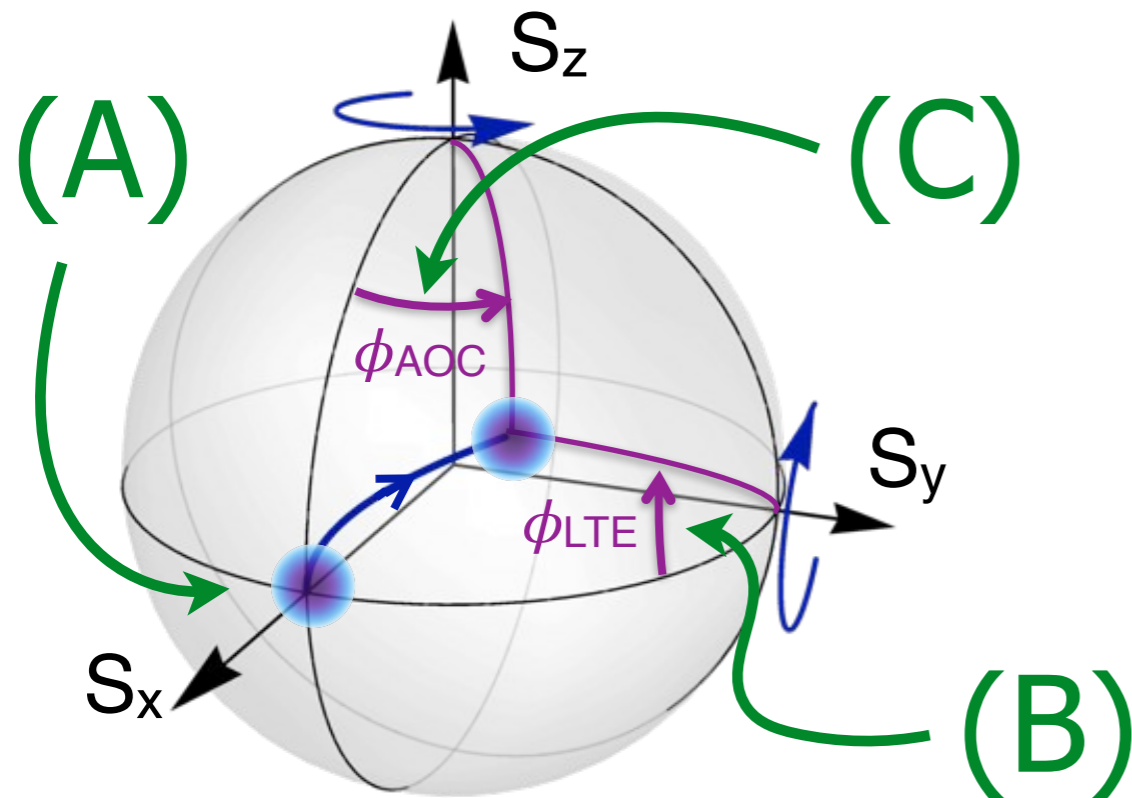
$$\Delta J_y = \frac{4}{G_2} \frac{1}{\sqrt{N}}$$

Alignment-to-orientation conversion



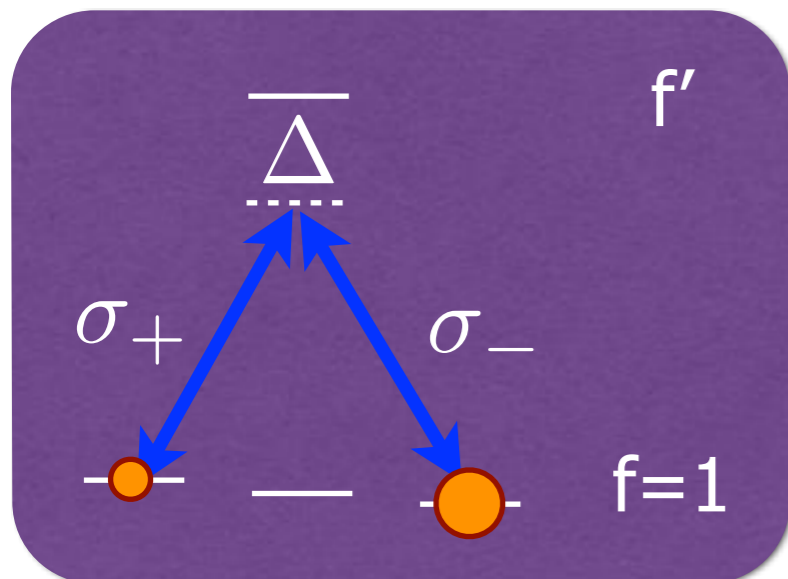
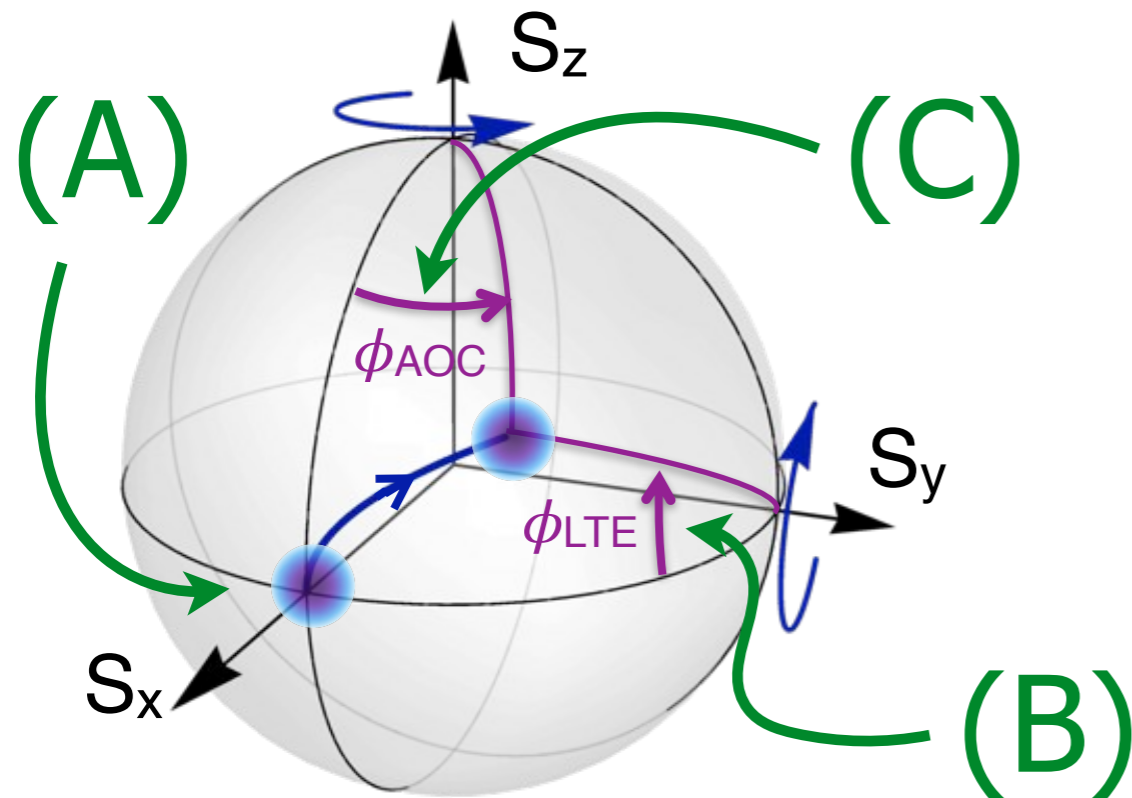
- (A) input $J_x = N_A/2$
aligned state
prepared via optical
pumping
- (B) B_z rotates collective
atomic spin
 $J_x \Rightarrow J_y \propto B_z$
- (C) probe pulses drive
Raman transition
 $J_y \Rightarrow J_z \propto S_x$

Nonlinear Faraday rotation



- (A) input $S_x = N_L/2$
polarised optical pulse
- (B) rotates $S_x \Rightarrow S_z \propto J_y$
(atomic birefringence)
- (C) rotates $S_x \Rightarrow S_y \propto J_z$
(Faraday rotation)

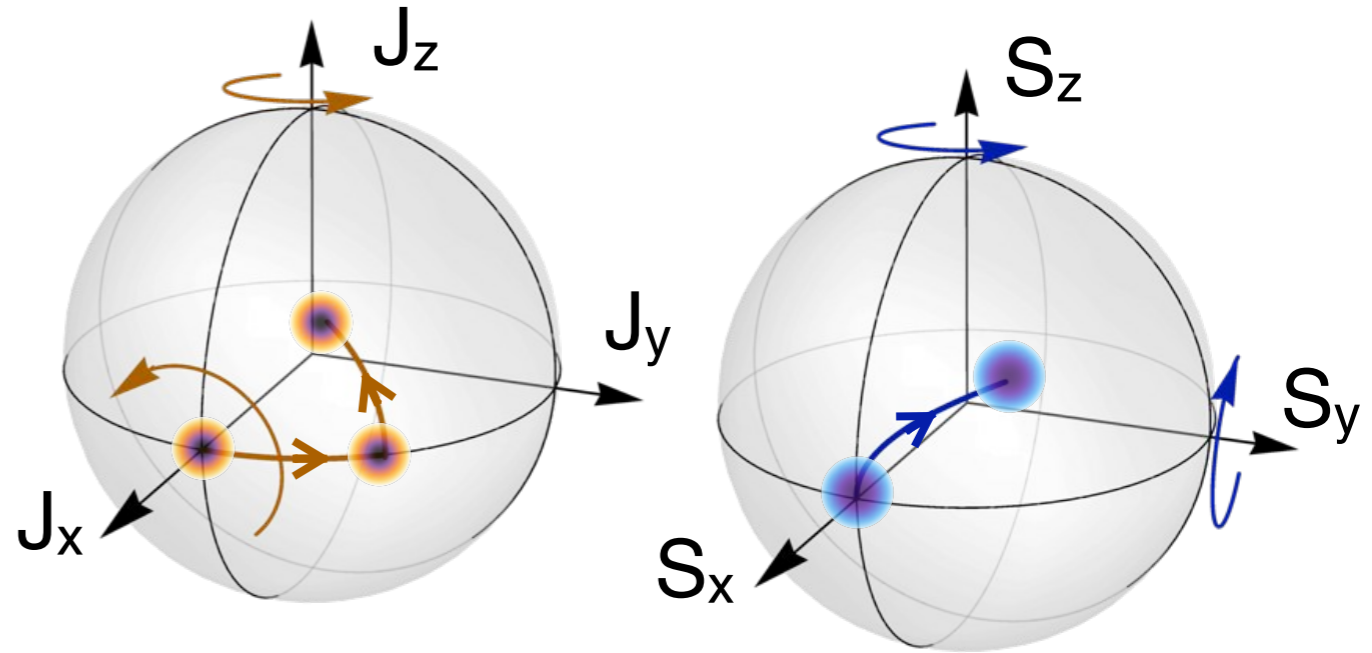
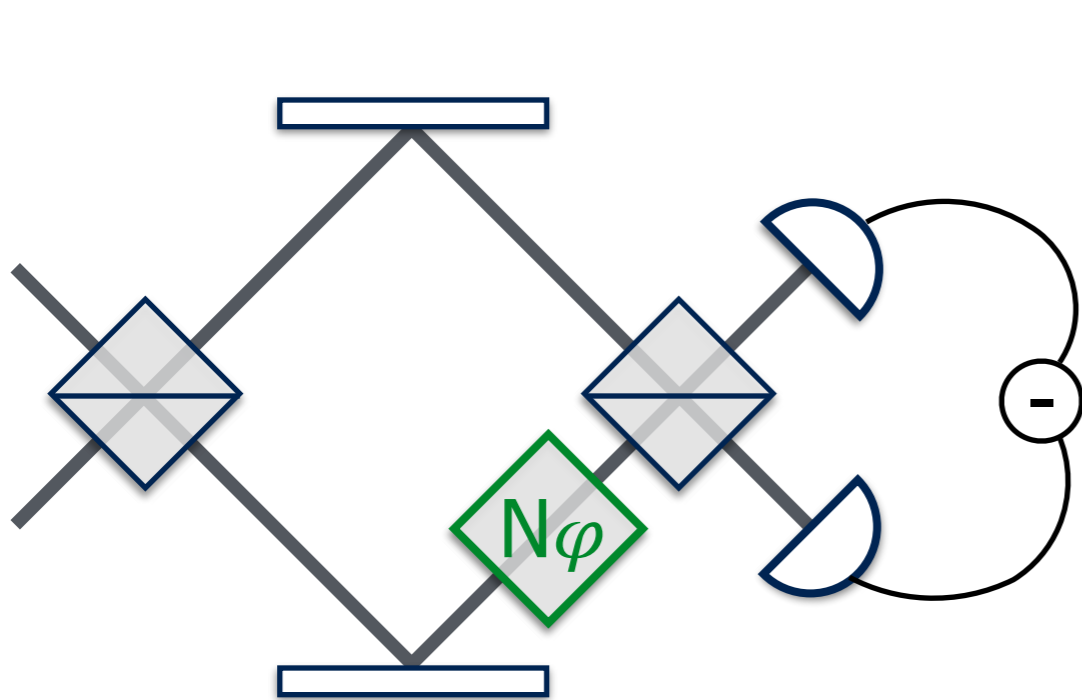
Nonlinear Faraday rotation



$$J_y \Rightarrow J_z \propto S_x$$

- (A) input $S_x = N_L/2$ polarised optical pulse
- (B) rotates $S_x \Rightarrow S_z \propto J_y$ (atomic birefringence)
- (C) rotates $S_x \Rightarrow S_y \propto J_z$ (Faraday rotation)

Nonlinear interferometer



phase shift $\phi = \frac{G_1 G_2}{4} N J_y$

shot-noise $\Delta\phi = 4/\sqrt{N}$

sensitivity

$$\Delta J_y = \frac{4}{G_1 G_2} \frac{1}{N^{3/2}}$$

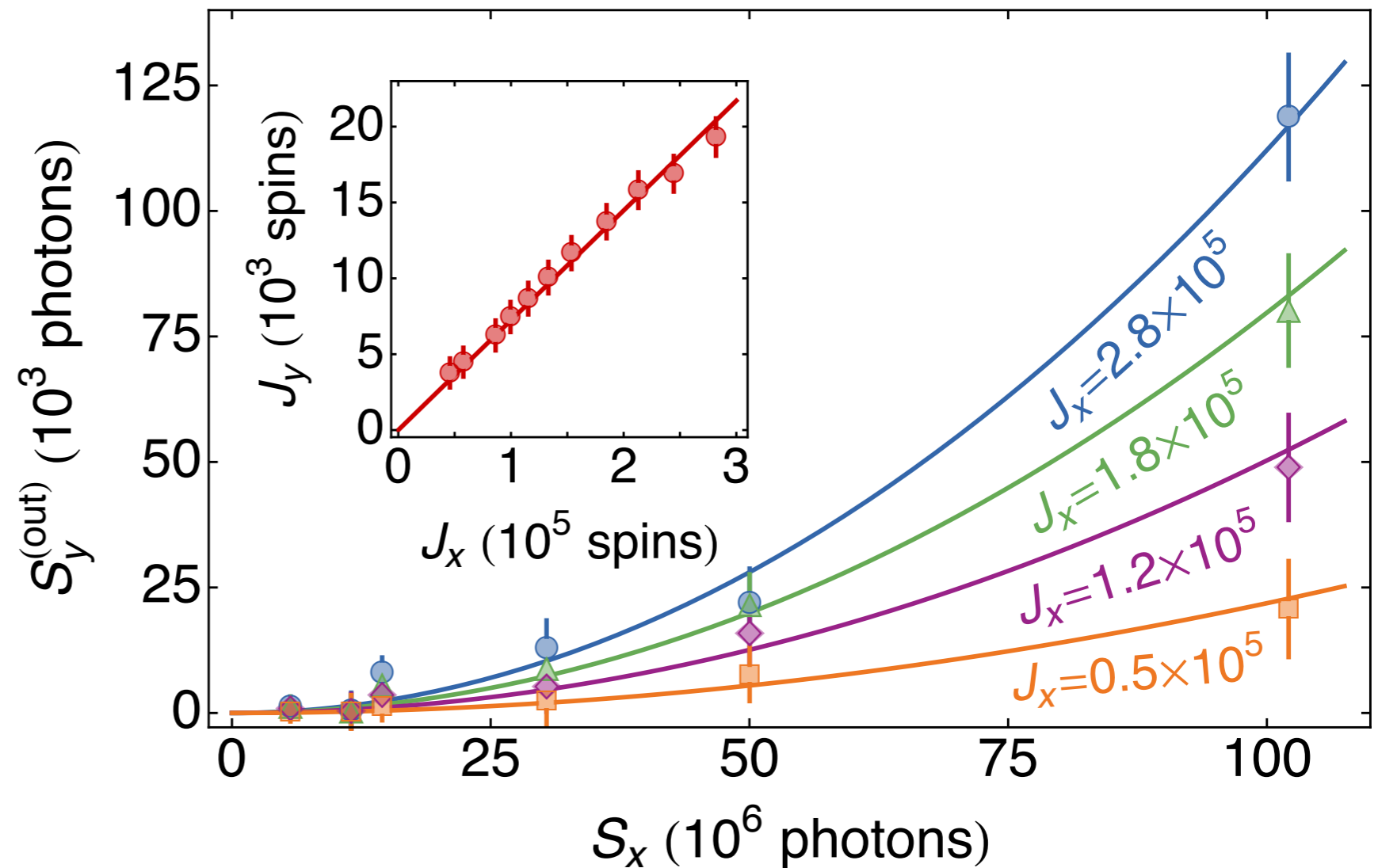
Nonlinear signal enhancement

signal from
nonlinear Faraday
rotation

$$S_y = \frac{G_1 G_2}{2} S_x^2 J_y$$

magnetometer
read-out

$$J_y \propto B_z$$



Measurement sensitivity

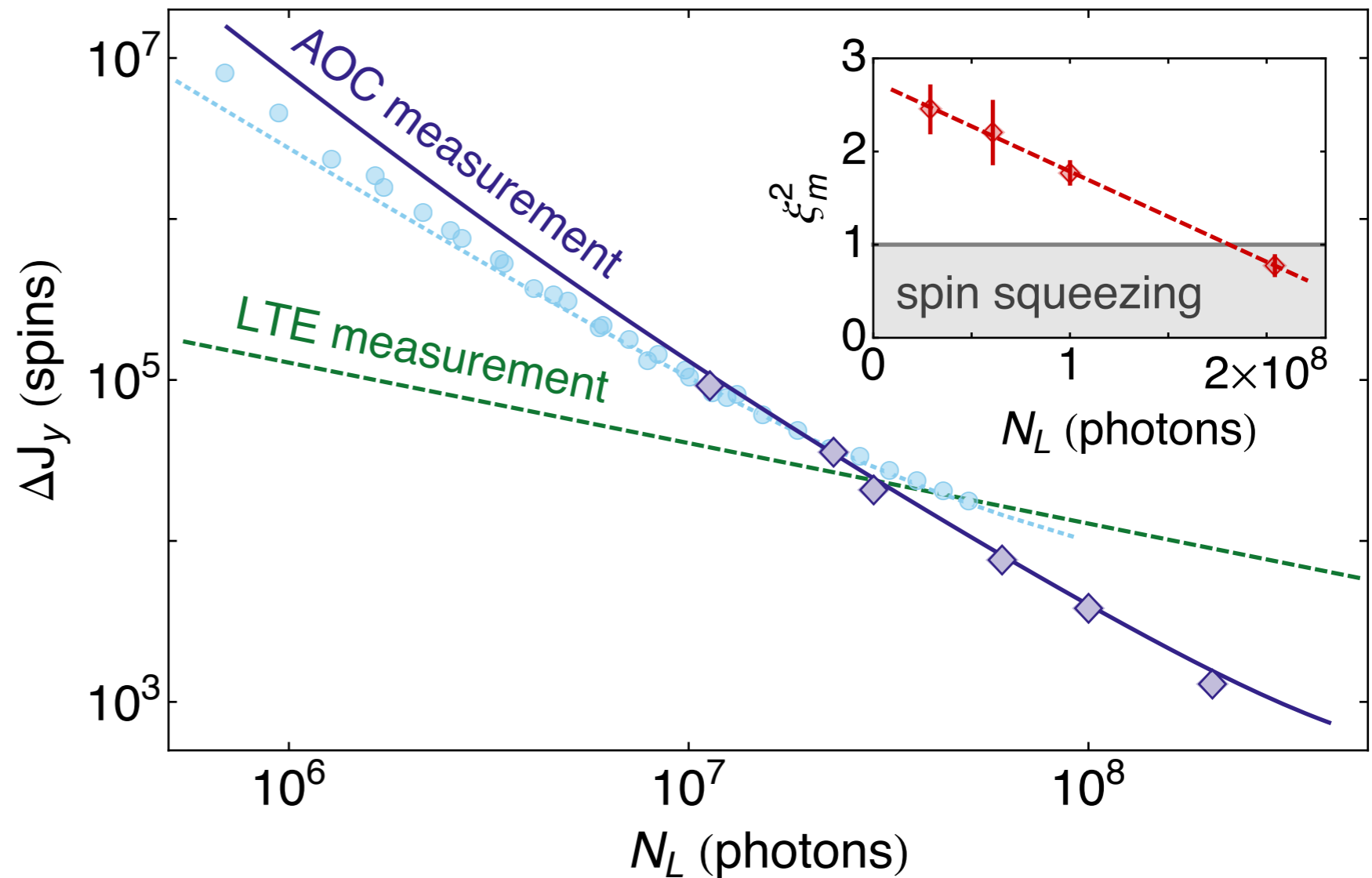
sensitivity

$$\Delta J_y = \frac{4}{G_1 G_2} \frac{1}{N^{3/2}}$$

$$\Delta J_y \simeq 10^3 \text{ spins}$$

spin squeezing

$$\xi^2 = \frac{(\Delta K_\theta)^2 J_x}{2 |J_x^{(\text{out})}|^2}$$
$$= 0.7$$



Measurement sensitivity

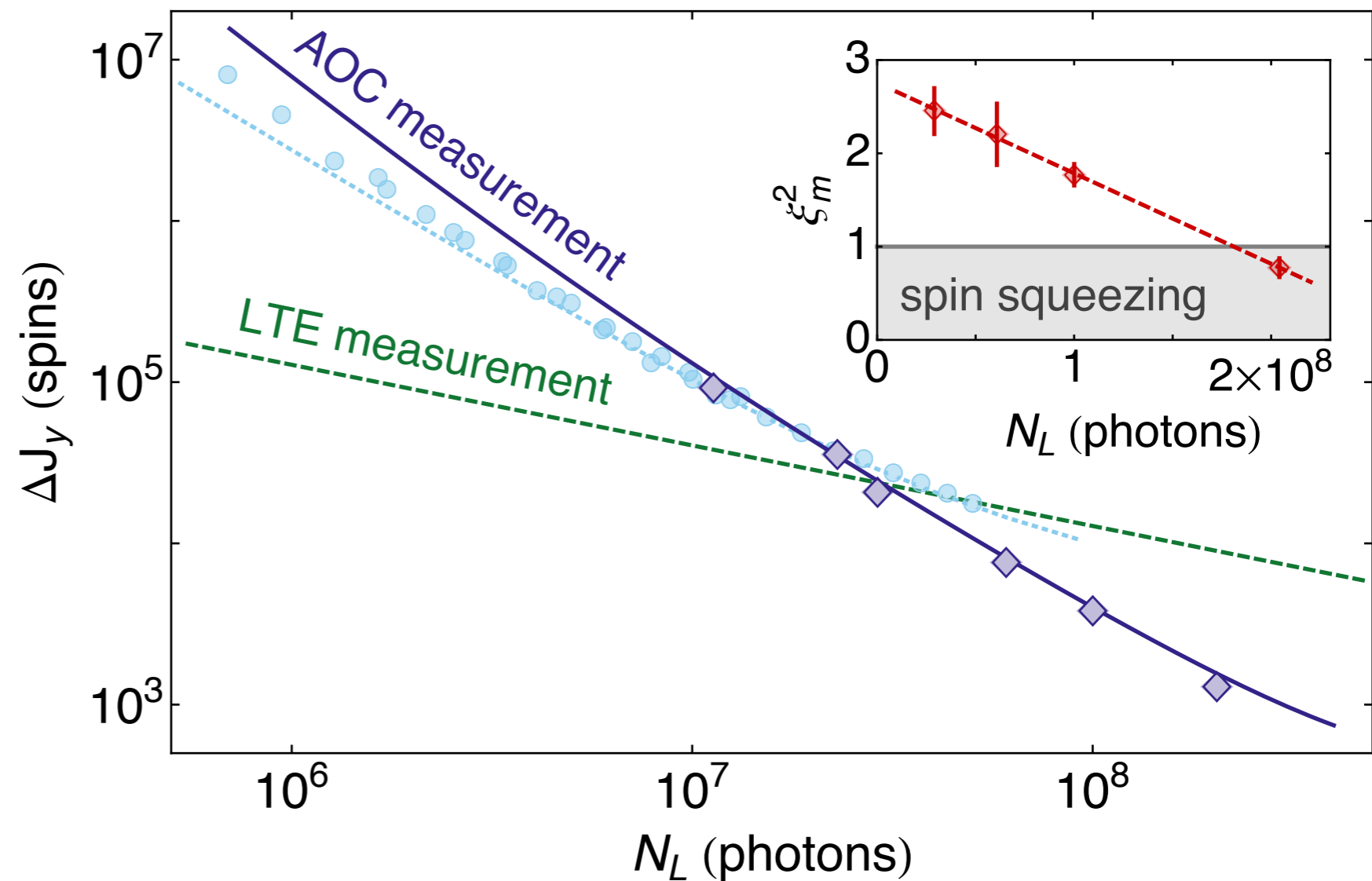
nonlinear read-out sensitivity

$$\Delta J_y = \frac{4}{G_1 G_2} \frac{1}{N^{3/2}}$$

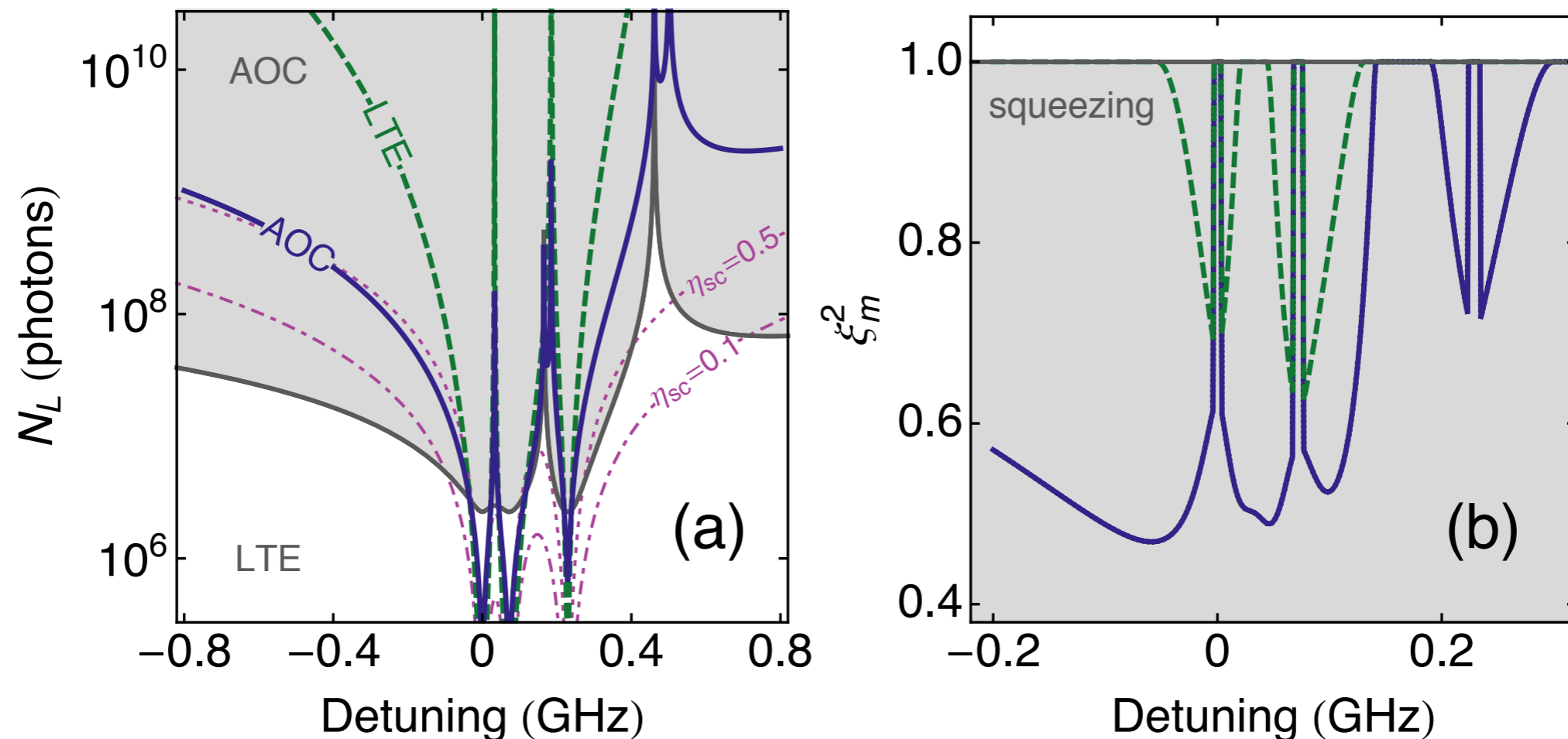
linear read-out sensitivity

$$\Delta J_y = \frac{4}{G_2} \frac{1}{\sqrt{N}}$$

crossover with
 $N=3 \times 10^7$ photons



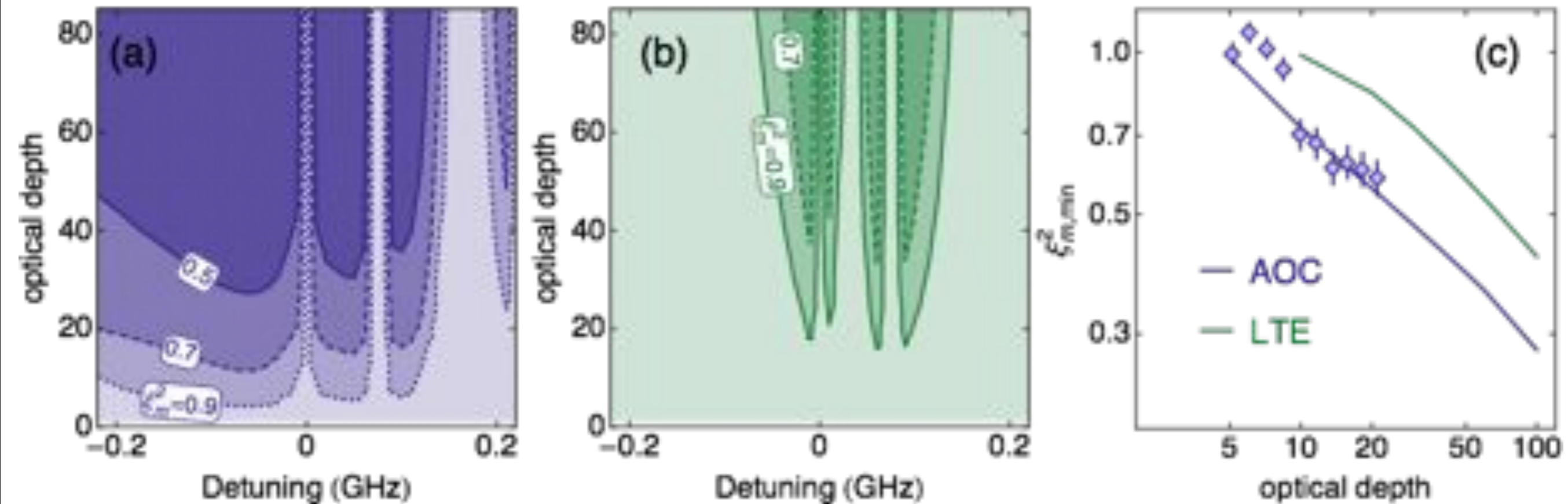
Nonlinear beats linear read-out



we optimise measurement sensitivity with respect to:

- probe detuning Δ
- probe power N (photon number)
- sample optical depth d_0 (atom number, interaction strength)

Nonlinear beats linear read-out



we optimise measurement sensitivity with respect to:

- probe detuning Δ
- probe power N (photon number)
- sample optical depth d_0 (atom number, interaction strength)

Can improved scaling give better sensitivity?

enhanced scaling

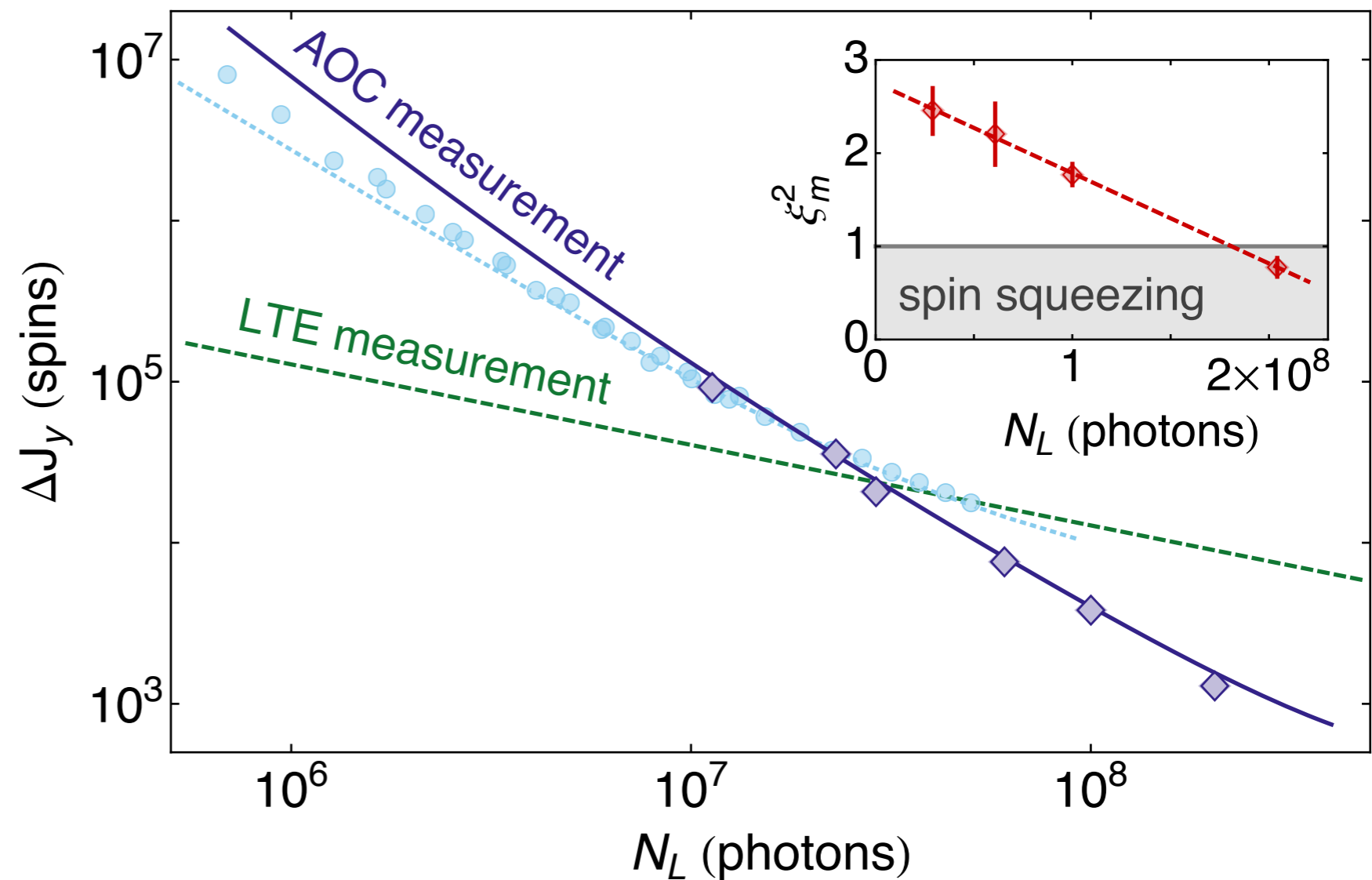
$$\Delta J_y \propto N^{-3/2}$$

metrologically
significant

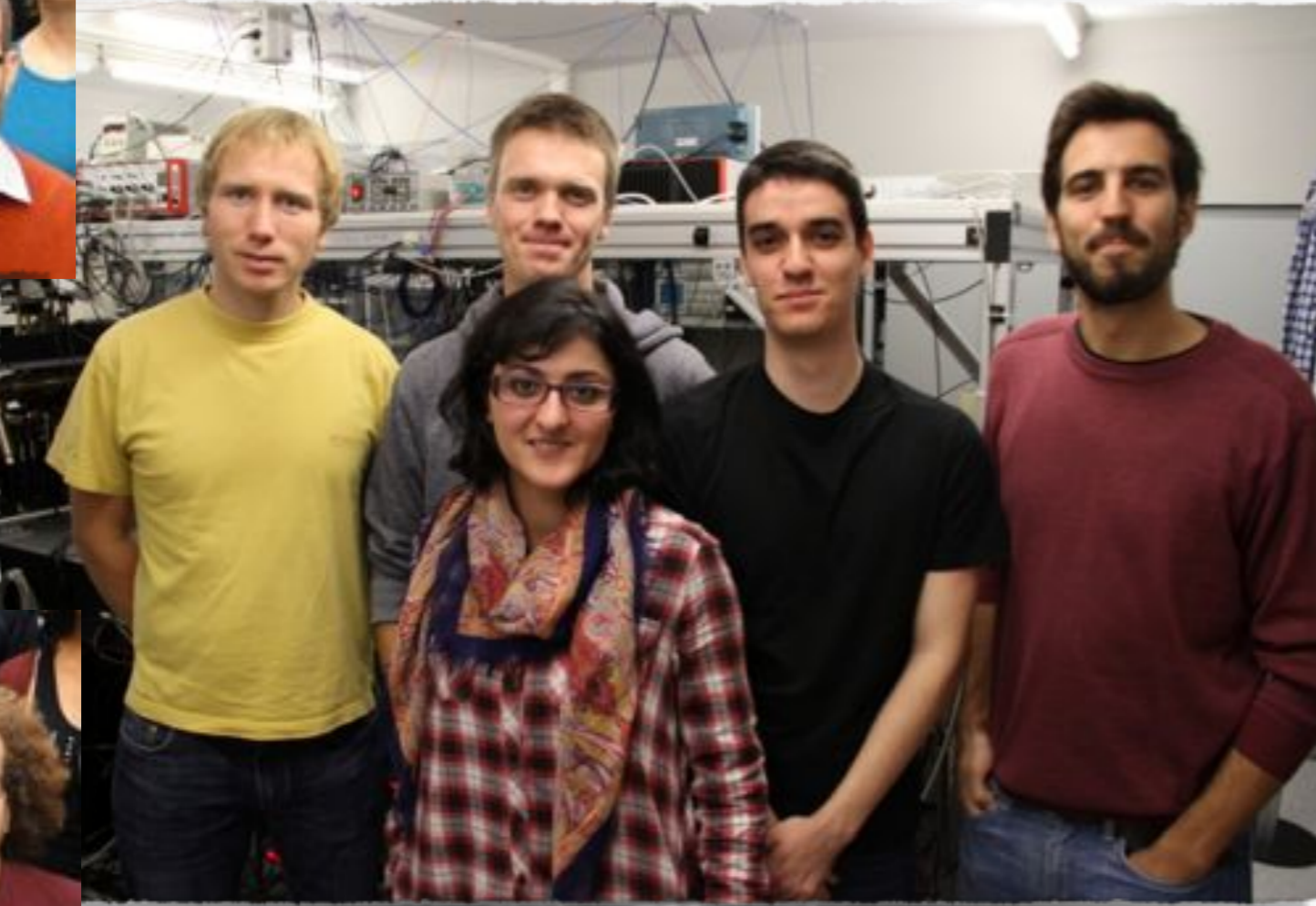
$$J_y \propto B_z$$

better absolute
sensitivity

vs. (N, Δ, d_0)



Acknowledgments



www.mitchellgroup.icfo.es

RJS et al. arXiv:
1310.5889 (2013)

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G. Colangelo &
M.W. Mitchell*

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